

CLASSICAL NUCLEATION AND LATTICE MODEL UNITE

ABSTRACT

Classical nucleation theory predicts the limit of superheat of liquids quite well. To come up with an equation for the limit of superheat of polymer solutions, the lattice model for polymer solutions was used to give the surface tension of polymer solutions. A formula for bubble nucleation in polymer solutions was derived by Jennings with the precursor equation $d \ln A / dK = 1 / (6K)$ where $J = A \exp K$ gives the nucleation rate for liquids. In this paper it is shown that the precursor equation holds for monomer in the polystyrene-cyclohexane system. Thus, the precursor equation is true for all molecular weight polymer.

NOMENCLATURES

- a surface area of solvent molecule
- A prefactor
- B coefficient
- C constant of integration
- d density of liquid
- J nucleation rate
- k Boltzmann constant
- K exponent
- M molecular weight of liquid
- MW_i molecular weight: solvent 1, polymer 2
- P_e equilibrium vapor pressure
- P_L ambient pressure
- P_v vapor pressure
- T temperature Kelvin
- T_c critical temperature in Kelvin
- w weight fraction polymer
- δ Poynting correction factor

ΔT rise in superheat in Centigrade

σ surface tension

INTRODUCTION

Classical nucleation theory goes back to Josiah Willard Gibbs where increasing entropy and decreasing energy give way to an emerging bubble forming aided by the rise in temperature known as the limit of superheat. Flory and Huggins theory provided an understanding of the behavior of a polymer solution and Prigogine and Marechal 1952 derived equations for the surface tension of polymer solutions.

THEORY

This paper is based on a highly cited paper by Blander and Katz 1975 and another paper by Siow and Patterson 1973. Jennings joined these streams of thought for the first time in 2012 by coming up with this formula for the limit of superheat as polymer $w_2 \rightarrow 0$.

$$\lim_{w_2 \rightarrow 0} \partial T / \partial w_2 = (MW_1 / MW_2) (3kT^2 / \sigma_1 a) \quad (1)$$

Mathematically, Eq. (1) is only true in the limit of zero weight fraction polymer in the solution because of the limitation in getting the surface tension. However, the actual Jennings and Middleman 1985 data shows straight vectors in T, w space for lower molecular weight. In Jennings' 2014 paper, the 2012 equation (1) was expanded because of that. Here is the expanded equation.

$$\Delta T = 3kT^2 w MW_1 / \sigma a MW_2 \quad (2)$$

So, the question is, What led to Eq. (2)? It is Eq. (3) from Jennings 2012 paper.

$$\lim_{w_2 \rightarrow 0} (\partial \ln A / \partial w_2) / (\partial K / \partial w_2) = 1 / (6K) \quad (3)$$

Another way to write (3) is (4).

$$d \ln A / dK = 1 / (6K) \quad (4)$$

A and K are found in the equation (Z) in the APPENDIX, Blander and Katz's formula for bubble nucleation

in liquids. A is the prefactor in (Z) and K is the exponent in (Z) . In the appendix is sketched out my method for checking the validity of (4) for monomer, that is $r = 1$, where r is the ratio of the molar volume of polymer to the molar volume of solvent. Furthermore, (4) has a well-defined integration constant that was calculated in Jennings 2020 article.

METHODOLOGY

In the early 1980s, Jennings and Middleman collected data on liquid→gas nucleation in polymer solutions that was published in 1985. In early 2020 the study was complete, except direct calculation revealed that the surface tension is so influential that only the surface tension terms matter.

RESULTS AND DISCUSSION

Unfortunately, there is no known data, except a lone styrene point in the 1985 paper, but plugging in the numbers gave a result that was off by about a factor of 10. There is data on polystyrene/benzene gathered by Prud'homme and Gregory 1985, but it does not agree with Eq. (2) at all. The beauty of this entire study from Fall 1980 to now is that it shows how the very successful classical nucleation theory unites with the lattice model for polymer solutions. The author knows of no one else who used Siow and Patterson's 1973 theory to solve this problem and the others who attempted it had complicated equations.

CONCLUSIONS

The author first started on this project in 1980 and feels that this paper concludes it. It is remarkable how influential the surface tension is in bubble nucleation. Possibly this paper will get noticed and lead others to explore the ramifications of the uniting of what started with Gibbs and Flory/Huggins.

APPENDIX

Part 1

Here, we establish $d \ln A / dK = 1 / (6K)$ for $r = 1$.

Proof of $d \ln A / dK = 1 / (6K)$ for all pure liquids $J = A \exp K$

from Blander and Katz classical nucleation theory liquid→gas.

The starting equations are from

1) Blander and Katz (Z) and 2) Siow and Patterson (A) and (B).

See JH Jennings, International Journal of Thermodynamics article, Ref. (7).

page 127-128. Everything here is done as $\varphi_2 \rightarrow 0$.

$$J \approx 3.73 \times 10^{35} (d^2 \sigma / M^3 B)^{1/2} \exp [- 1.182 \times 10^5 \sigma^3 / (T (P_V - P_L)^2)] \quad (Z)$$

$$(\sigma - \sigma_1) a / kT = \ln (\varphi_{1S} / \varphi_1) + ((r - 1) / r) (\varphi_{2S} - \varphi_2) \quad (A)$$

$$\ln [(\varphi_{2S} / \varphi_2)^{1/r} / (\varphi_{1S} / \varphi_1)] = (\sigma_1 - \sigma_2) a / kT \quad (B)$$

Now, near $\varphi_2 = 0$, Eq. (B) becomes

$$\varphi_{2S} = \varphi_2 \exp [r (\sigma_1 - \sigma_2) a / kT]$$

Putting in the numbers, $\partial \varphi_{2S} / \partial \varphi_2 \approx 10^{-38}$ for $MW_2 = 2000$, $r = 13.4$ and

even less for higher MW. The object here is to let $r = 1$ near $\varphi_2 = 0$.

Thus $\varphi_{2S} = \varphi_2 \exp [(\sigma_1 - \sigma_2) a / kT]$ and so $(\partial \varphi_{2S} / \partial \varphi_2) = \exp [(\sigma_1 - \sigma_2) a / kT]$.

All of the details had to be included because it was only evident later that the surface tension term dominated. This had to be shown by direct calculation.

$$1. d = d_1 + (d_2 - d_1) \varphi_{2S}$$

$$\partial d / \partial \varphi_2 = (d_2 - d_1) (\partial \varphi_{2S} / \partial \varphi_2) = (d_2 - d_1) \exp [(\sigma_1 - \sigma_2) a / kT]$$

$$2. Pe = Pe(0) \varphi_{1S} = Pe(0) (1 - \varphi_{2S})$$

$$\partial Pe / \partial \varphi_2 = - Pe(0) (\partial \varphi_{2S} / \partial \varphi_2) = - Pe(0) \exp [(\sigma_1 - \sigma_2) a / kT]$$

3. $B \approx 1 - 1/3 (1 - P_L / P_V)$ Here the approximation $P_e = P_V$ is used because

this is a very small correction and they are close. Hence,

$$\begin{aligned} \partial B / \partial \phi_2 &= (-1/3) (-1) \partial / \partial \phi_2 P_L / P_e \\ &= (1/3) P_L \partial / \partial \phi_2 1 / P_e \text{ and } = (1/3) P_L / P_e(0) \exp [(\sigma_1 - \sigma_2) a / kT] \end{aligned}$$

$$4. \delta = 1 - d_G/d + 0.5 (d_G/d)^2 \text{ and } d_G = P_e MW_1/RT_1$$

$$d_G/d = x \text{ is small so } \delta \approx e^{-x}$$

After some manipulation we have.

$$\partial \delta / \partial \phi_2 = \delta [d_G/d + (d_G/d)^2 (d_2 - d_1)] \exp [(\sigma_1 - \sigma_2) a / kT]$$

Note ideal gas and T_1 and ϕ_2 are orthogonal. Also note that these differentials don't vanish.

It was seen that the surface tension influence dominates wherever it occurs.

$\partial \phi_{2s} / \partial \phi_2 \approx$ only vanishes with r greater than 5, say.

To avoid clutter, the values for all these unknown are left out, but they are all in Jennings 2012

International Journal of Thermodynamics article. This part is to establish the method.

Now, we get $\lim_{\phi_2 \rightarrow 0} \partial \sigma / \partial \phi_2 = (k T_o / a) (1 - \exp [(\sigma_1 - \sigma_2) a / kT]) = (k T_o / a) (0.9984)$.

This is arrived at by letting $r = 1$ in (A) above and partial differentiating with respect to ϕ_2 .

Part 2

Getting all these quantities is the hard part and the rest is calculating

$\lim_{\phi_2 \rightarrow 0} (\partial \ln A / \partial \phi_2) / (\partial K / \partial \phi_2) = 1 / (6K)$, where the A and K are the prefactor/exponent in (Z) above.

For $r = 1$, $1 / (6K) = -0.002399$, or $K = -69.47$.

From the IAJER article by Jennings 2020, $r = 13.4$ and $K = -64.56$ for 10 solvents, polar & non-polar.

Deriving Jennings 2012 Equation (1) in the THEORY section depends on $|K|$ being large. The surface tension

dominates, as is noted by Blander and Katz 1975. The validity of $d\ln A/dK = 1/(6K)$ is therefore true for all r

and its integration constant for non-polar and polar solvents was calculated in Jennings' IAJER

2020 article. All of this puts Jennings' 2014 Polymers Research Journal equation on solid ground, which is

formula (2) in THEORY, the equation for bubble nucleation in polymer solutions. The polymer solution data

for this was first published by Jennings/Middleman 1985 MACROMOLECULES.

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